

#### Precision Absolute Luminosity with Photon Pairs

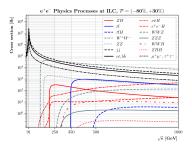
Graham W. Wilson

University of Kansas

October 11, 2023

Using the pure QED process  $e^+e^- \to \gamma\gamma$  for luminosity measurements.

# LUMI: Targets for Absolute Luminosity Precision



- The standard process used for **absolute** luminosity at LEP is small-angle **Bhabha** scattering,  $e^+e^- \rightarrow e^+e^-$  (high statistics).
- This will be important for **relative** luminosity and could still lead in absolute precision.
- The pure QED process,  $e^+e^- \rightarrow \gamma\gamma$ , is now also considered very seriously for **absolute** luminosity, for both experimental and theoretical reasons.
- It emphasizes reconstruction (rejection) of high energy photons (electrons) over most of the detector's solid angle.
- Ideally match/exceed stat. precision of the accelerator. Denominator normalizing processes should have cross-sections exceeding the numerator.
- Example 1 (ILC): WW at 250 GeV. With 0.9  $\mathrm{ab}^{-1}$  (LR)  $\rightarrow$  1.7  $\times$  10^{-4}.
- Example 2 (10^{12} Z with FCC)  $\rightarrow 1.0 \times 10^{-6}.$

What is realistically achievable in terms of systematics is another matter. For now the assumption is to target  $10^{-4}$ .

### LUMI: $e^+e^- \rightarrow \gamma\gamma$ for absolute luminosity

Targeting  $10^{-4}$  precision. Cross-sections (and ratios) at  $\sqrt{s} = 161$  GeV.

$\theta_{\min}$ (°)	$\sigma_{\gamma\gamma}$ (pb)	$\Delta\sigma/\sigma$ (10 $\mu$ rad)	$\sigma(ee)/\sigma(\gamma\gamma)$
45	5.3	$2.0 imes10^{-5}$	6.1
20	12.7	$2.2 imes10^{-5}$	22
15	15.5	$2.4 imes10^{-5}$	35
10	19.5	$2.9 imes10^{-5}$	68
6	24.6	$3.9 imes10^{-5}$	155
2	35.7	$8.1 imes10^{-5}$	974

Unpolarized Born cross-sections. ±24% for (80%/30%) longitudinal beam polarization. Typical HO effects: + 5 to 10%. Counting statistics adequate for √s ≫ m<sub>Z</sub>. Note: Use whole detector.

• For comparison, 10 $\mu$ rad knowledge for OPAL small-angle **Bhabha** lumi acceptance, corresponds to uncertainty of 100  $\times$  10<sup>-5</sup>.

 $\gamma\gamma$  has "relaxed" fiducial acceptance tolerances compared to Bhabhas.

• Bhabha rejection (e/ $\gamma$  discrimination) important. Can be aided by much better azimuthal measurements given electron bending in the B-field. FoM: *B* z<sub>LCAL</sub>. ILD has 7.7 Tm. FCC about 2.2 Tm. OPAL was 1.04 Tm. Adequate rejection feasible within tracker acceptance? / challenging below.

# Why is $e^+e^- \rightarrow \gamma \gamma$ attractive?

Focus here on experimental things. The hope and expectation is that theory will be able to keep up.

- Bhabha process looks problematic for precision absolute luminosity. It was even not under control experimentally at LEP1 due to the beam-beam effect bias on the luminosity acceptance at the 0.1% level (see 1908.01704).
- Di-photon process should be less affected.
- Di-photons much less sensitive to polar angle metrology than Bhabhas.
- Di-photons less sensitive to FSR than Bhabhas.
- Likely more feasible now with modern calorimeters to do a particle-by-particle reconstruction. Likely easier with di-photons.
- Current detector designs are arguably over-designed for Bhabhas with some compromises for overall performance especially for high energy photons in azimuthal and energy reconstruction, and perhaps for hermeticity.
- Di-photons at very low angle is challenging! but gives significant added value to the assumed clean measurements in the tracker acceptance.

So let's design precision forward calorimetry for electrons AND photons inspired by various ideas (and avoiding some of the compromises) of related designs, CALICE, ILD, SiD, CMS-HGCAL, ALICE-FoCal, Fermi-LAT.

# PLUG-Cal: Precision Luminosity Ultra-Granular Calo.

#### Initial Design Ideas

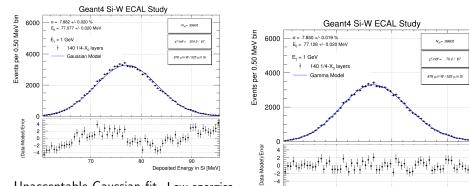
- Precise location of the high-energy photon interaction point (via conversion to  $e^+e^-$ ) in thin absorbers (see Fermi-LAT for extreme version of this).
- 250 GeV photons need longitudinal containment to avoid large constant term. (10, 1)% of photons survive for (3, 6) X<sub>0</sub> prior to interaction.
- $\textbf{0} \text{ Above items} \rightarrow \text{Many thin layers assuming a sampling Si-W ECAL}.$
- Calibration  $\rightarrow$  More straightforward with uniform sampling.
- Potential for adoption in part of pixel-based devices. FoCal prototype achieved 30 micron resolution for high energy electron showers with ALPIDE sensors (1708.05164). 2 planes adopted for ALICE-FoCal upgrade.
- **(**) Include  $0^{\text{th}}$ -layer and maybe more for enhanced  $e/\gamma$  discrimination.
- Emphasize azimuthal measurements for  $e^+e^- / \gamma \gamma$  discrimination. Expect about 110 mrad acoplanarity for  $Bz_{LCAL} = 8.4$  Tm.
- O Particle-by-particle reconstruction capabilities.
- $\textbf{O} \ \ \text{Limited solid-angle} \rightarrow \text{cost is not an over-arching concern.}$

### PLUG-Cal: Initial GEANT4 Design Studies

- In collaboration with Brendon Madison. We have been exploring some aspects of the design using various GEANT4 (4-11-01-patch-02 [MT]) examples (TestEm3, HGCAL\_testbeam, gammaray\_telescope)
- Pasic EM energy performance studies using TestEm3. Range cut 1 micron. XY extent 50 cm. Adds up globally the energies deposited in each type of material. Apply to Si-W calorimeter with various absorber and sensor thicknesses. Main results are for 35 X<sub>0</sub> depth of W absorber with 140 samples with same Si sensor thickness as ILD.
- Also recently (Saturday...) started with HGCAL\_testbeam example looking at position resolution observables. This has hexagonal pads with similar transverse dimensions to standard ILD and SiD.

### Measuring Energy Linearity and Resolution

Typical calorimeter analyses fit Gaussian distributions to truncated regions of plots. Here instead a Gamma distribution is used to also model the skewness. The **two** parameters can be configured to be the mean,  $\mu$ , and the fractional resolution,  $(\sigma'/\mu)$ . The mean and fractional resolution are annotated as  $(E_0, \sigma)$  in the plots.



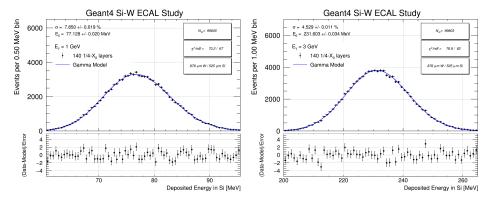
Unacceptable Gaussian fit. Low energies and worse designs give distinct positive skew. Not surprising given what we know about the Poisson and Landau distributions.

But fits great to Gamma. As  $\sigma$  improves it tends to a Gaussian (CLT).

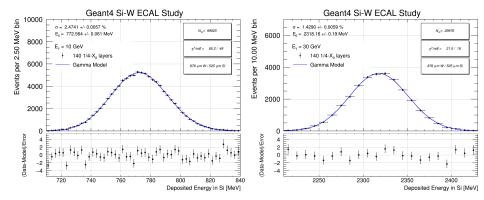
80

Deposited Energy in Si [MeV

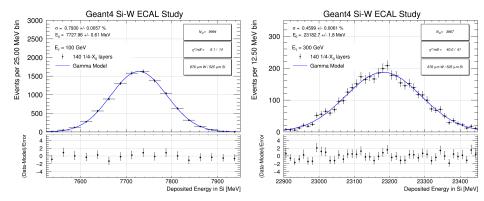
### Energy Linearity and Resolution: 1 GeV, 3 GeV Photons



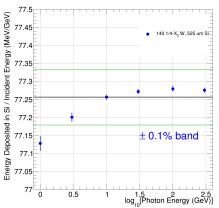
### Energy Linearity and Resolution: 10 GeV, 30 GeV Photons



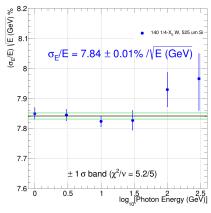
### Energy Linearity and Resolution: 100, 300 GeV Photons



# Energy Linearity and Resolution



Calorimeter Photon Linearity



Calorimeter Photon Energy Resolution

Excellent linearity in [1, 300] GeV range. Generally within 0.1%, suspect albedo for < 2 GeV. EM sampling fraction of 7.7%.

Fits well with only a stochastic term and **no** constant term. Energy resolution of  $0.460 \pm 0.006\%$  at 300 GeV.

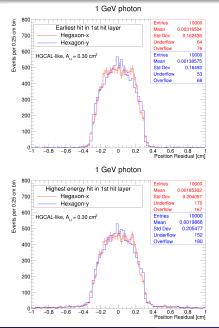
#### Position Resolution Tests

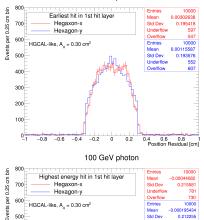
How much can the photon and electron position resolution be pushed with small cells? Can one localize the initial photon interaction point? thus measuring the  $\gamma$  scattering angle,  $\theta = \tan^{-1}(r/z)$ , and aiding in separating electrons and photons.

- Use GEANT4 example HGCal\_testbeam (CMS). The software was well adapted to the task but is NOT the proposed design concept.
- Uses **hexagonal** Si pads with 28 layers totalling 27 X<sub>0</sub>. Absorbers included Pb, Cu, CuW (quite a mix...).
- In a first step changed hexagonal pixel areas from  $1.09 \text{ cm}^2$  to  $0.301 \text{ cm}^2$ .
- So far, longitudinal structure unchanged except beam starts inside Al box. Beam particles are incident on the array with a Gaussian profile with spread in x and y of 1.5 cm. Residuals for calorimeter position observables are calculated with respect to the randomized true beam position event-by-event.



#### Choosing the best hit in the first hit layer





Hexagon-y

HGCAL-like, A. = 0.30 cm<sup>2</sup>

600

500

400

300

200

100

-0.8

-0.6 -0.4 -0.2 781

730

737

813

10000

0.212255

-0.000195434

Underflow

Overflow

Entries

Mean

Std Dev

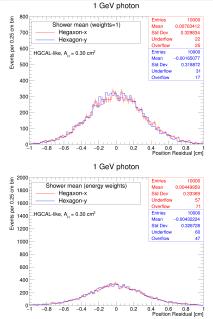
Underflow

Overflow

0.6

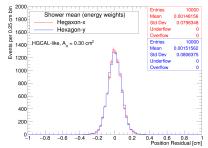
<sup>100</sup> GeV photon

### Shower center-of-gravity (all layers)



#### h 800 Entries 10000 Shower mean (weights=1) Mean 0.00280931 Events per 0.25 cm 700 Hegaxon-x Std Dev 0.250091 Underflow Hexagon-y Overflow 14 600 HGCAL-like, A. = 0.30 cm<sup>2</sup> Entries 10000 Mean 0.00336975 500 Std Dev 0.252861 Underflow Overflow 15 400 300 200 100 -0.8 -0.6 -0.4 -0.20.4 0.6 0.8 Position Residual [cm]

#### 100 GeV photon

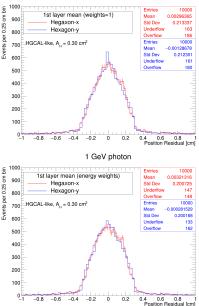


Graham W. Wilson (University of Kansas)

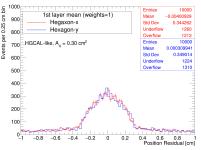
100 GeV photon

### First Hit Layer CoG

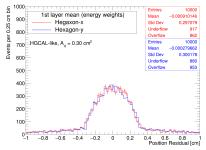
1 GeV photon



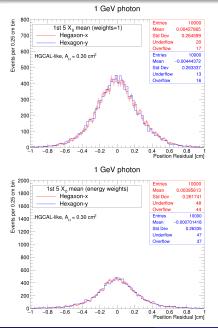
#### 100 GeV photon

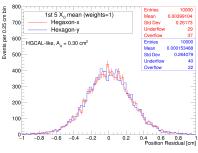


#### 100 GeV photon

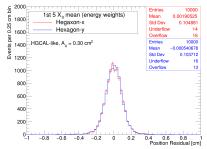


#### CoG from layers within 5 $X_0$ of 1st hit layer





#### 100 GeV photon



#### 100 GeV photon

- Good sensitivity at the single cell level for low energy photons.
- More ambiguities for higher energy photons, but much more information from whole shower.
- Much higher granularity can benefit a lot. See eg. FoCal prototype. Dimensions (in microns) of 50\*50, 30\*30, 25\*100, 12.5\*50 are all possibilities for pure digital approach.
- Need to also make sure that layer-to-layer alignment is randomized enough.
- Need to do some clustering too.
- Hexagons are different!

- I believe the PLUG-Cal concept has potential for superior performance for luminosity measurements even with  $e^+e^- \rightarrow \gamma\gamma$  below the tracker acceptance. Potential doubling of acceptance.
- It can likely make radial measurements better than ILD LumiCal but with longer Moliere radius and better energy and azimuthal resolutions and hermeticity.
- Plan to benchmark against current ILD design for electrons and photons.
- What fraction if any of digital-only planes not clear. Could also consider analog + digital planes if digital thin enough. I'm wary of compromising the analog performance.